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ON THE "MOVEMENT" OF THE ZEROS OF EIGENFUNCTIONS OF THE STURM-LIOUVILLE PROBLEM

TIGRAN HARUTYUNYAN, AVETIK PAHLEVANYAN, YURI ASHRAFIYAN

ABSTRACT. We study the dependence of the zeros of eigenfunctions of Sturm-Liouville problem on the parameters that define the boundary conditions. As a corollary, we obtain Sturm oscillation theorem, which states that the n -th eigenfunction has n zeros.

Let us consider Sturm-Liouville boundary value problem $L(q, \alpha, \beta)$:

$$ly \equiv -y'' + q(x)y = \mu y, \quad 0 < x < \pi, \quad \mu \in \mathbb{C}, \quad (1)$$

$$y(0) \cos \alpha + y'(0) \sin \alpha = 0, \quad \alpha \in (0, \pi], \quad (2)$$

$$y(\pi) \cos \beta + y'(\pi) \sin \beta = 0, \quad \beta \in [0, \pi), \quad (3)$$

where $q \in L^1_{\mathbb{R}}[0, \pi]$, i.e. q is a real-valued, summable function on $[0, \pi]$.

By $L(q, \alpha, \beta)$ we also denote the self-adjoint operator corresponding to the problem (1)–(3).

It is well-known, that the problem $L(q, \alpha, \beta)$ has a countable set of simple, real eigenvalues (see, e.g. [1, 2, 3, 4]), which we denote by $\mu_n(q, \alpha, \beta)$, $n = 0, 1, \dots$, (emphasizing the dependence on q , α and β) and enumerate in increasing order:

$$\mu_0(q, \alpha, \beta) < \mu_1(q, \alpha, \beta) < \dots < \mu_n(q, \alpha, \beta) < \dots \quad (4)$$

In the papers [4, 7] it was introduced the concept of the eigenvalues function (EVF) of a family of Sturm-Liouville operators $\{L(q, \alpha, \beta) : \alpha \in (0, \pi], \beta \in [0, \pi)\}$. For fixed q it is a function of two variables γ and δ ($\gamma = \alpha + \pi n \in (0, \infty)$, $\delta = \beta - \pi m \in (-\infty, \pi)$) determined through the eigenvalues $\mu_n(q, \alpha, \beta)$, $n = 0, 1, 2, \dots$, by the following formula:

$$\mu(\gamma, \delta) = \mu(\alpha + \pi n, \beta - \pi m) := \mu_{n+m}(q, \alpha, \beta), \quad n, m = 0, 1, 2, \dots \quad (5)$$

It is proved that this function is analytic with respect to γ and δ . It is strictly increasing by γ and strictly decreasing by δ .

It is also known (see, e.g. [1]) that every nontrivial solution $y(x, \mu)$ of the equation (1) may have only simple zeros (if $y(x_0, \mu) = 0$, then $y'(x_0, \mu) \neq 0$), and (see, e.g. [5]) every solution $y(x, \mu)$ is a continuously differentiable function with respect to the totality of variables x and μ . Therefore, by applying the implicit function theorem (see, e.g. [6, p. 452]), we get that the zeros of the solution $y(x, \mu)$ are the continuously differentiable functions with respect to μ . Since the function $x = x(\mu)$, such that the identity $y(x(\mu), \mu) \equiv 0$ is true for all μ from some interval

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(a, b) , is called the solution of the equation $y(x, \mu) = 0$, then differentiating the last identity with respect to μ we obtain:

$$\frac{dy(x(\mu), \mu)}{d\mu} = \frac{\partial y(x(\mu), \mu)}{\partial x} \frac{dx(\mu)}{d\mu} + \frac{\partial y(x(\mu), \mu)}{\partial \mu} \equiv 0, \mu \in (a, b). \quad (6)$$

Let us denote $\frac{\partial y(x, \mu)}{\partial \mu} := \dot{y}(x, \mu)$ and write the identity (6) in the following form:

$$\frac{dx(\mu)}{d\mu} = \dot{x}(\mu) = -\frac{\dot{y}(x(\mu), \mu)}{y'(x(\mu), \mu)}, \mu \in (a, b). \quad (7)$$

On the other hand, let us write down the fact that $y(x, \mu)$ is the solution of the equation (1), i.e.

$$-y''(x, \mu) + q(x)y(x, \mu) \equiv \mu y(x, \mu), \quad 0 < x < \pi, \mu \in \mathbb{C}, \quad (8)$$

and differentiating this identity with respect to μ we receive:

$$-\dot{y}''(x, \mu) + q(x)\dot{y}(x, \mu) \equiv y(x, \mu) + \mu\dot{y}(x, \mu). \quad (9)$$

Multiplying (8) by $\dot{y}$, (9) by y and subtracting from the second obtained identity the first one, we get

$$y''(x, \mu)\dot{y}(x, \mu) - \dot{y}''(x, \mu)y(x, \mu) \equiv y^2(x, \mu), \quad 0 < x < \pi, \mu \in \mathbb{C},$$

i.e.

$$\frac{d}{dx}[y'(x, \mu)\dot{y}(x, \mu) - \dot{y}'(x, \mu)y(x, \mu)] \equiv y^2(x, \mu). \quad (10)$$

If we integrate this identity with respect to x from 0 to a ($0 \leq a \leq \pi$), then

$$\begin{aligned} y'(a, \mu)\dot{y}(a, \mu) - \dot{y}'(a, \mu)y(a, \mu) - y'(0, \mu)\dot{y}(0, \mu) + \\ + \dot{y}'(0, \mu)y(0, \mu) = \int_0^a y^2(x, \mu) dx, \end{aligned} \quad (11)$$

and if we integrate with respect to x from a to π , then we get

$$\begin{aligned} y'(\pi, \mu)\dot{y}(\pi, \mu) - \dot{y}'(\pi, \mu)y(\pi, \mu) - y'(a, \mu)\dot{y}(a, \mu) + \\ + \dot{y}'(a, \mu)y(a, \mu) = \int_a^\pi y^2(x, \mu) dx. \end{aligned} \quad (12)$$

Now, as $y(x, \mu)$ let us take $y = \varphi(x, \mu, \alpha, q)$ – the solution of the equation (1), satisfying the following initial conditions:

$$\varphi(0, \mu, \alpha, q) = \sin \alpha, \quad \varphi'(0, \mu, \alpha, q) = -\cos \alpha. \quad (13)$$

It is easy to see that eigenfunctions of the problem $L(q, \alpha, \beta)$ are obtained from the solution $\varphi(x, \mu, \alpha, q)$ at $\mu = \mu_n(q, \alpha, \beta)$ (here we use (5)), i.e.

$$\begin{aligned} \varphi_n(x, q, \alpha, \beta) &:= \varphi_n(x) = \varphi(x, \mu_n(q, \alpha, \beta), \alpha, q) = \\ &= \varphi(x, \mu(\alpha + \pi n, \beta), \alpha, q) = \varphi(x, \mu(\alpha, \beta - \pi n), \alpha, q) = \\ &= \varphi(x, \mu(\alpha, \delta), \alpha, q)|_{\delta=\beta-\pi n} = \varphi(x, \mu(\alpha, \delta))|_{\delta=\beta-\pi n} = \\ &= \varphi(x, \mu(\alpha, \beta - \pi n)). \end{aligned} \quad (14)$$

Let $0 \leq x_n^0 < x_n^1 < \dots < x_n^m \leq \pi$ be the zeros of the eigenfunction $\varphi_n(x, q, \alpha, \beta) = \varphi(x, \mu(\alpha, \beta - \pi n))$, i.e. $\varphi_n(x_n^k, q, \alpha, \beta) = \varphi(x_n^k, \mu(\alpha, \beta - \pi n)) = 0, k = 0, 1, \dots, m$.

Let q, α, n be fixed. We consider the following questions:

- a) How are the zeros $x_n^k = x_n^k(\beta)$, $k = 0, 1, \dots, m$ changing, when β is changing on $[0, \pi)$?
- b) How many zeros the n -th eigenfunction $\varphi_n(x, q, \alpha, \beta)$ has, i.e. what is equal m ?

By taking $y = \varphi_n(x)$ and $a = x_n^k$ ($k = 0, 1, \dots, m$) in (11), we receive

$$\begin{aligned} \varphi_n'(x_n^k) \dot{\varphi}_n(x_n^k) - \dot{\varphi}_n'(x_n^k) \varphi_n(x_n^k) - \varphi_n'(0) \dot{\varphi}_n(0) + \\ + \dot{\varphi}_n'(0) \varphi_n(0) = \int_0^{x_n^k} \varphi_n^2(x) dx. \end{aligned} \quad (15)$$

Since the initial conditions (13) should be held for all $\mu \in \mathbb{C}$, then $\dot{\varphi}_n(0) = 0$ and $\dot{\varphi}_n'(0) = 0$. Also, taking into account that $\varphi_n(x_n^k) = 0$, from (15) we get

$$\varphi_n'(x_n^k) \dot{\varphi}_n(x_n^k) = \int_0^{x_n^k} \varphi_n^2(x) dx. \quad (16)$$

Since the zeros of the solutions are simple, then $\varphi_n'(x_n^k) \neq 0$, and therefore, from (16) the following equality implies :

$$\frac{\dot{\varphi}_n(x_n^k)}{\varphi_n'(x_n^k)} = \frac{1}{(\varphi_n'(x_n^k))^2} \int_0^{x_n^k} \varphi_n^2(x) dx. \quad (17)$$

Now, from (7), we obtain

$$\dot{x}_n^k(\mu_n) = \left. \frac{dx_n^k(\mu)}{d\mu} \right|_{\mu=\mu_n} = -\frac{\dot{\varphi}_n(x_n^k)}{\varphi_n'(x_n^k)} = -\frac{1}{(\varphi_n'(x_n^k))^2} \int_0^{x_n^k} \varphi_n^2(x) dx, \quad (18)$$

i.e. zeros $x_n^k(\mu_n)$, $k = 0, 1, \dots, m$ of the eigenfunction $\varphi_n(x)$ are decreasing if the eigenvalue $\mu_n(q, \alpha, \beta)$ is increasing, which in its turn means that

$$\dot{x}_n^k(\mu_n(q, \alpha, \beta)) \leq 0. \quad (19)$$

Let us note, that the equality $\dot{x}_n^k(\mu_n) = 0$ may occur only at $x_n^k = 0$, i.e., when $x = 0$ is a zero of the eigenfunction $\varphi_n(x)$, and it is so at $\alpha = \pi$ ($\gamma = \pi l$, $l = 1, 2, \dots$).

Meanwhile, in the inequality $\dot{x}_n^k(\mu_n) = \dot{x}_n^k(\mu_n(q, \alpha, \beta)) \leq 0$, the variable μ_n can be changed depending on q , α and β . More precisely, if for some change in these three variables $\mu_n(q, \alpha, \beta)$ increases, then the zeros of the eigenfunction $\varphi_n(x)$ are moving to the left, and if $\mu_n(q, \alpha, \beta)$ decreases, then the zeros of the eigenfunction $\varphi_n(x)$ are moving to the right.

In the work [8], it was proved for $q \in L_{\mathbb{R}}^2[0, \pi]$ that the number of zeros of the n -th eigenfunctions of the problems $L(q, \pi, 0)$ and $L(0, \pi, 0)$ are equal. But it is easy to see that the same proof remains true for $q \in L_{\mathbb{R}}^1[0, \pi]$.

It is easy to calculate that the eigenvalues of the problem $L(0, \pi, 0)$ are $\mu_n(0, \pi, 0) = (n+1)^2$ and the eigenfunctions are

$$\begin{aligned} \varphi_n(x) &= \varphi(x, \mu_n(0, \pi, 0), \pi, 0) = \\ &= \varphi\left(x, (n+1)^2, \pi, 0\right) = \frac{\sin(n+1)x}{n+1}, \quad n = 0, 1, 2, \dots \end{aligned} \quad (20)$$

The zeros of this eigenfunction are $x_n^k = \frac{\pi k}{n+1}$, $k = 0, 1, \dots, n+1$, i.e. the n -th eigenfunction of the problem $L(0, \pi, 0)$ has $n+2$ zeros in $[0, \pi]$, two of which are 0 and π , and n zeros in $(0, \pi)$.

Thereby, the n -th eigenfunction $\varphi(x, q, \pi, 0)$ of the problem $L(q, \pi, 0)$ has two zeros at the endpoints of $[0, \pi]$, i.e. $x_n^0(q, \pi, 0) = 0$, $x_n^{n+1}(q, \pi, 0) = \pi$ and also n zeros in the interval $(0, \pi)$.

Along with increasing β from 0 to π the eigenvalue $\mu_n(q, \pi, 0)$ is continuously (with respect to β) decreasing from $\mu_n(q, \pi, 0)$ to $\mu_n(q, \pi, \pi) = \mu(\pi, \pi - \pi n) = \mu(\pi, 0 - (n-1)\pi) = \mu_{n-1}(q, \pi, 0)$ (see [5]) and, according to (18), the zeros of the function $\varphi_n(x, q, \pi, \beta)$ are increasing, i.e. are moving to the right (all but leftmost zero $x_n^0 = 0$). In particular, the rightmost zero $x_n^{n+1} = \pi$, by moving to the right, leaves the segment $[0, \pi]$ and in $[0, \pi]$ remains $n+1$ zeros (one $x_n^0 = 0$ and n zeros in the interval $(0, \pi)$). And the previous zero reaches π when the relation $\varphi_n(\pi) = \varphi(\pi, \mu_n(q, \pi, \beta), \pi, q) = c_n \psi_n(\pi) \sin \beta = 0$ again occurs (see below (21) and (24)), and this is possible only when β reaches π (and $\mu_n(q, \pi, \beta)$, by decreasing, reaches $\mu_n(q, \pi, \pi) = \mu_{n-1}(q, \pi, 0)$). Then the eigenfunction $\varphi_n(x) = \varphi(x, \mu_n(q, \pi, \beta), \beta, q)$ will smoothly transform to the eigenfunction $\varphi(x, \mu_n(q, \pi, \pi), \pi, q) = \varphi(x, \mu_{n-1}(q, \pi, 0), \pi, q)$ which has $n+1$ zeros in $[0, \pi]$, two out of which are the endpoints 0 and π , and $n-1$ zeros are in $(0, \pi)$. Thus, the oscillation theorem is proved for all $L(q, \pi, \beta)$, $\beta \in [0, \pi]$.

Now, as $y(x, \mu)$ let us take $\psi(x, \mu, \beta, q)$ – the solution of the equation (1) satisfying the following initial conditions:

$$\psi(\pi, \mu, \beta, q) = \sin \beta, \quad \psi'(\pi, \mu, \beta, q) = -\cos \beta. \quad (21)$$

It is easy to see that the eigenvalues $\mu_n = \mu_n(q, \alpha, \beta)$, $n = 0, 1, \dots$, of the problem $L(q, \alpha, \beta)$ are the zeros of the entire function

$$\Psi(\mu) = \Psi(\mu, \alpha, \beta, q) = \psi(0, \mu, \beta, q) \cos \alpha + \psi'(\pi, \mu, \beta, q) \sin \alpha, \quad (22)$$

and the eigenfunctions, corresponding to these eigenvalues, are obtained by the formula

$$\psi_n(x) = \psi(x, \mu_n(q, \alpha, \beta), \beta, q), \quad n = 0, 1, \dots \quad (23)$$

Since all eigenvalues μ_n are simple, then eigenfunctions $\varphi_n(x)$ and $\psi_n(x)$ corresponding to the same eigenvalue μ_n are linearly dependent, i.e. there exist the constants $c_n = c_n(q, \alpha, \beta)$, $n = 0, 1, \dots$, such that

$$\varphi_n(x) = c_n \psi_n(x), \quad n = 0, 1, \dots \quad (24)$$

This implies that $\varphi_n(x, q, \alpha, \beta)$ and $\psi_n(x, q, \alpha, \beta)$ have equal number of zeros.

Let $0 \leq x_n^m < x_n^{m-1} < \dots < x_n^0 \leq \pi$ be the zeros of the eigenfunction $\psi_n(x, q, \alpha, \beta)$ i.e. $\psi_n(x_n^k, q, \alpha, \beta) = 0$, $k = 0, 1, \dots, m$.

By taking in the identity (12) $y = \psi_n(x) = \psi(x, \mu_n(q, \alpha, \beta), \beta, q)$, we receive

$$\begin{aligned} \psi_n'(\pi) \dot{\psi}_n(\pi) - \dot{\psi}_n'(\pi) \psi_n(\pi) - \psi_n'(x_n^k) \dot{\psi}_n(x_n^k) + \\ + \dot{\psi}_n'(x_n^k) \psi_n(x_n^k) = \int_{x_n^k}^{\pi} \psi_n^2(x) dx. \end{aligned} \quad (25)$$

From (21) we get that $\dot{\psi}_n(\pi) = 0$ and $\dot{\psi}_n'(\pi) = 0$. And since $\psi_n(x_n^k) = 0$, then the equality (25) gets the form

$$-\psi_n'(x_n^k) \dot{\psi}_n(x_n^k) = \int_{x_n^k}^{\pi} \psi_n^2(x) dx. \quad (26)$$

As all zeros x_n^k are simple, i.e. $\psi'_n(x_n^k) \neq 0$, dividing both sides of the last equality by $(\psi'_n(x_n^k))^2$, we obtain

$$\frac{\dot{\psi}_n(x_n^k)}{\psi'_n(x_n^k)} = -\frac{1}{(\psi'_n(x_n^k))^2} \int_{x_n^k}^{\pi} \psi_n^2(x) dx. \quad (27)$$

Now from (7), by taking $y = \psi_n(x)$, $x = x_n^k$, we have

$$\dot{x}_n^k(\mu_n) = \left. \frac{dx_n^k(\mu)}{d\mu} \right|_{\mu=\mu_n} = -\frac{\dot{\psi}_n(x_n^k)}{\psi'_n(x_n^k)} = \frac{1}{(\psi'_n(x_n^k))^2} \int_{x_n^k}^{\pi} \psi_n^2(x) dx \geq 0, \quad (28)$$

i.e. the zeros $x_n^k(\mu_n)$, $k = 0, 1, \dots, m$ of the eigenfunction $\psi_n(x)$ are increasing, if the eigenvalue μ_n is increasing. Note that the equality $\dot{x}_n^k(\mu_n) = 0$ is possible only when $x_n^0(\mu_n) = \pi$, but it holds when $\beta = 0$ ($\delta = -\pi l$, $l = 0, 1, 2, \dots$).

While studying the dependence of the zeros of eigenfunctions on α , it is convenient to use formula (28), because the eigenfunctions $\psi_n(x)$ have fixed values $\psi_n(\pi, \mu, \beta) = \sin \beta$ and $\psi'_n(\pi, \mu, \beta) = -\cos \beta$ for all $\mu \in \mathbb{C}$, i.e. all $\psi_n(x)$ satisfy the initial conditions $\psi_n(\pi) = \sin \beta$ and $\psi'_n(\pi) = -\cos \beta$, which means that from the endpoint π of the segment $[0, \pi]$ (when changing α) new zeros can neither enter nor leave (neither appear nor disappear). Thus, with increasing α , the eigenvalues $\mu_n(q, \alpha, \beta)$ (with fixed q and β) are increasing and according to (28) (i.e. $\dot{x}_n^k(\mu_n) \geq 0$) the zeros of the eigenfunction $\psi_n(x)$ are moving to the right (i.e. are increasing). Wherein, the values $\psi(\pi) = \sin \beta$ and $\psi'(\pi) = -\cos \beta$ are not changed, the number of the zeros increase, and these zeros can neither "collide" nor "split" as they are simple. That's why new zeros can appear only entering through 0– left endpoint of the segment $[0, \pi]$ and moving to the right (and "condensing" respectively). And new zeros are entered through 0– left endpoint of the segment $[0, \pi]$ only when $\psi_n(0) = 0$, but since $\psi_n(0) = c_n \varphi_n(0) = c_n \sin \alpha$ ($c_n \neq 0$), then the equality $\psi_n(0) = 0$ is possible only when $\sin \alpha = 0$, i.e. in our notations only when $\alpha = \pi$ (and because $\mu_n(q, 0, \beta) = \mu(0 + \pi n, \beta) = \mu(\pi + (n-1)\pi, \beta) = \mu_{n-1}(q, \pi, \beta)$) or when $\alpha = 0$.

Hence, when $\alpha = \pi$, the eigenfunction $\psi_n(x, q, \pi, \beta)$ as well as $\varphi_n(x, q, \pi, \beta)$ has n zeros in $(0, \pi)$ and one zero at $x = 0$ left endpoint. Wherein, $\psi_n(x, q, \pi, \beta) = \psi(x, \mu_n(q, \pi, \beta), \beta, q) = \psi(x, \mu_{n+1}(q, 0, \beta), \beta, q) = \psi_{n+1}(x, q, 0, \beta)$. With increasing α from 0 to π the eigenvalue $\mu_{n+1}(q, 0, \beta)$ is increasing (continuously with respect to α) to $\mu_{n+1}(q, \pi, \beta)$, and the leftmost zero $x = 0$ by moving to the right, appears in $(0, \pi)$, i.e. there are $n+1$ zeros of the eigenfunction $\psi_n(x, q, \alpha, \beta)$ in $(0, \pi)$ (and another fixed zero $x = \pi$, if $\beta = 0$). A new zero will appear at $x = 0$ left endpoint, when α will reach $\alpha = \pi$.

Thus, we obtain the following oscillation theorem:

Theorem 1. *The eigenfunctions of the problem $L(q, \alpha, \beta)$ corresponding to the n -th eigenvalue $\mu_n(q, \alpha, \beta)$, $n = 0, 1, 2, \dots$, have exactly n zeros in $(0, \pi)$. All these zeros are simple. If $\alpha = \pi$ and $\beta = 0$, then the n -th eigenfunction has $n+2$ zeros in $[0, \pi]$, and if either $\alpha = \pi$, $\beta \in (0, \pi)$ or $\beta = 0$, $\alpha \in (0, \pi)$, then the n -th eigenfunction has $n+1$ zeros in $[0, \pi]$.*

Oscillation properties of the solutions of the problem $L(q, \alpha, \beta)$ (Sturm theory), the studies of which initiated by Sturm in [9, 10], outlined in monographic literature (see, e.g. [1, 11, 12]) for continuous q . In the recent years, the study was mostly focused on the cases when q is bounded or $q \in L^2_{\mathbb{R}}[0, \pi]$ (see, e.g. [8, 13]), but

in many studies (see [14] and references therein) implicitly assumes that Sturm's oscillation theorem (that the n -th eigenfunction has n zeros) is also true for $q \in L^1_{\mathbb{R}}[0, \pi]$, although the rigorous proof is not available in the literature.

Our oscillation theorem is true for all $q \in L^1_{\mathbb{R}}[0, \pi]$.

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TIGRAN HARUTYUNYAN

FACULTY OF MATHEMATICS AND MECHANICS, YEREVAN STATE UNIVERSITY, 1 ALEX MANOOGIAN, 0025, YEREVAN, ARMENIA

E-mail address: hartigr@yahoo.co.uk

AVETIK PAHLEVANYAN

FACULTY OF MATHEMATICS AND MECHANICS, YEREVAN STATE UNIVERSITY, 1 ALEX MANOOGIAN, 0025, YEREVAN, ARMENIA

E-mail address: avetikpahlevanyan@gmail.com

YURI ASHRAFYAN

FACULTY OF MATHEMATICS AND MECHANICS, YEREVAN STATE UNIVERSITY, 1 ALEX MANOOGIAN, 0025, YEREVAN, ARMENIA

E-mail address: yuriashrafyan@ysu.am